

**Outdoor Sound Propagation with
Analytic Ray Curve Tracer and Gaussian Beam**

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Abstract

Outdoor sound propagation benefits from algorithms that can handle, in a computationally efficient manner, inhomogeneous media, complex boundary surfaces, and large spatial expanse. One recent work⁴⁵ proposed a ray tracing method using analytic ray curves as tracing primitives, which improved the performance of propagation paths computation over rectilinear ray tracers. In this paper, an algorithm is developed that extends the performance improvement to field computation; it combines the analytic ray curve tracer⁴⁵ with fast pressure computation based on the Gaussian beam model. The algorithm is validated against published results on benchmarks in atmospheric and ocean acoustics, and its application is demonstrated on a scene with terrains and buildings of realistic complexity and under a variety of atmospheric conditions. This algorithm is able to compute characteristic sound fields for fully general media profiles and complex 3D scenes at close-to-interactive speed.

I. INTRODUCTION

Sound propagation in outdoor environments^{1;2;3}, including atmospheric and underwater acoustics, must deal with spatially varying as well as moving media. The media profiles used as input to sound propagation can come from empirical models, real-world measurements, or from computational flow simulation. Obstacles' shape and material properties play an important role as well, especially for scenes with complex terrains or area with dense man-made structures. As data that describe the media conditions and the scene obstacles become increasingly available with ever richer details, methods for outdoor sound propagation needs to be able to account for the full scale of those data.

Existing methods face many challenges in handling such complexity; they either make assumptions that preclude a fully general media or complex obstacles in the scene, or they become prohibitively expensive with large, general scenes. Models such as Fast Field Program (FFP), Parabolic Equation (PE), and normal modes fall in the former category, while methods like Finite Difference Time Domain (FDTD), Finite Elements Method (FEM), and Boundary Element Method (BEM) belong to the latter.

Geometric acoustics (GA)⁵ methods like ray models are known for their efficiency in handling boundary surfaces under the assumption of homogeneous media and rectilinear paths. Recent works^{44;45} has also attempted to accommodate inhomogeneous media by tracing parabolic ray curves as primitives, which significantly accelerates path

computation. On the other hand, ray models suffer from artifacts in caustics and shadow zones when computing fields; while models such as the Gaussian beam⁶ perform better in this regard. However, the performance of Gaussian beam can be hindered by the underlying numerical path integration, which remains slow for inhomogeneous media.

We combine the performance of the analytic ray curve tracer⁴⁵ and the accuracy of the Gaussian beam⁶ into an algorithm for outdoor sound propagation. In particular:

1. We compute analytic solutions to on-ray pressure as well as near-ray fields based on the parabolic ray formulation (Sec. III), which leads to efficient field computation that matches the efficiency of the path computation.
2. We combine the Gaussian beam model with the analytic ray tracer and validate the approach on 2D benchmarks^{42;43} that are widely used in atmospheric and ocean acoustics. Our algorithm is able to replicate the published reference results generated by alternative techniques (Sec. IV).
3. We apply the algorithm on a 3D scene consisting of thousands of surface primitives modeling terrains and buildings for a set of different atmospheric conditions, which demonstrates its efficiency in computing characteristic sound fields (Sec. V).

Overall, we provide a validated solution to outdoor sound propagation that augments a fast analytic ray tracer with equally fast analytic field computations. This algorithm takes

general media and scene input and computes the full 3D sound field at close-to-interactive speed, making it useful for a wide range of outdoor sound applications (Sec. VI).

II. PRIOR WORK

Outdoor sound propagation has been studied extensively in underwater³ and atmospheric acoustics¹. Here we first review numerical methods that provide full wave solutions, including hybrid schemes aiming at reducing the high computation cost of wave-based methods. Next, we review ray-based methods and their advantages and limitations for outdoor scenarios.

A. Wave-based Methods

Early methods, including Fast Field Program (FFP) and Parabolic Equation (PE), provide frequency-domain, full-wave solutions that account for the inhomogeneous media, but depend upon simplifying assumptions about the media and scene configurations. In scenarios that meet those assumptions, these models have been thoroughly validated^{7;8;9;10} and often serve as reference solutions to test other models.

Among the general numerical methods that handle arbitrary media and obstacles, one widely-used method is the Finite Difference Time Domain (FDTD) method, which solves the linearized Euler equation^{11;12}. FDTD has been coupled with flow simulation^{13;14}; it has been applied with various ground conditions¹⁵, terrains^{16;17;18}, and complex

obstacles^{19;20;21}. As a time-domain method, it is also used for pulse propagation^{22;20}. The main disadvantage of FDTD is its limited scalability with domain volume or frequency; this limitation makes it prohibitively expensive for large-scale broadband simulation, and limits its usefulness for wide area to low-frequency cases²³. Methods such as Pseudo Spectral Time Domain (PSTD) and Transmission Line Matrix (TLM) are more efficient, but they are still fundamentally limited by the cost of discretizing a large domain.

To address the scalability problem, many hybrid methods have been developed. Some use FDTD in confined areas and apply PE to propagation over long range and relatively sparse space^{24;25;26}. Others, such as BEM or the Equivalent Source Method (ESM), were employed to limit the computation either to boundary surfaces or to volumes that bound scatterers tightly. But these methods introduced new issues. BEM, which must be coupled with specialized Green function for refractive media, scales poorly with surface area and frequency. ESM, when coupled with ray models²⁷ to handle large domains, does not scale well with the number or complexity of scatterer objects. A more recent method, Adaptive Rectangular Decomposition (ARD)^{28;29}, took advantage of the analytic solution of the wave equation in a rectangular domain, but it requires constant sound speed within each spatial subdivision, which is not easily adapted to a general media profile.

B. Geometric Acoustics Methods

Geometric acoustics (GA) methods⁵ are widely used in room acoustics⁴ to handle the

high-order surface interactions under the valid assumption of a homogeneous media.

Examples of GA methods include the image source method^{30;31}, ray tracing^{32;33}, beam tracing³⁴, and path tracing³⁵.

Ray models have also been applied to inhomogeneous media³ by numerically integrating the ray equations. While a sparse set of rays can be efficiently traced to plot out the propagation paths, long-range propagation and pressure field computation that requires dense rays become expensive. When the ray models are used to compute the pressure field, they are known to have issues in the caustic zones and in the shadow zones. The Gaussian beam approach⁶, which was developed in seismology and applied on underwater³⁷ and atmospheric³⁶ acoustics, improves the accuracy in caustics and shadow zones. However, when the underlying paths are still computed by numerical ray integration, the performance is limited by the integration step sizes. One recent work⁴⁵ achieved significant performance improvement by replacing the numerical ray integration with segments of parabolic ray curves. We give an overview of this ray tracer in the next section, and then introduce our algorithm, which combines this ray tracer with the Gaussian beam.

III. ALGORITHM

An overview of our algorithm is illustrated in Figure 1. Our algorithm is built upon an efficient ray tracer that outputs propagation paths made up of segments of analytic curves (Sec. A). We compute a set of additional variables for each ray segment by analytic

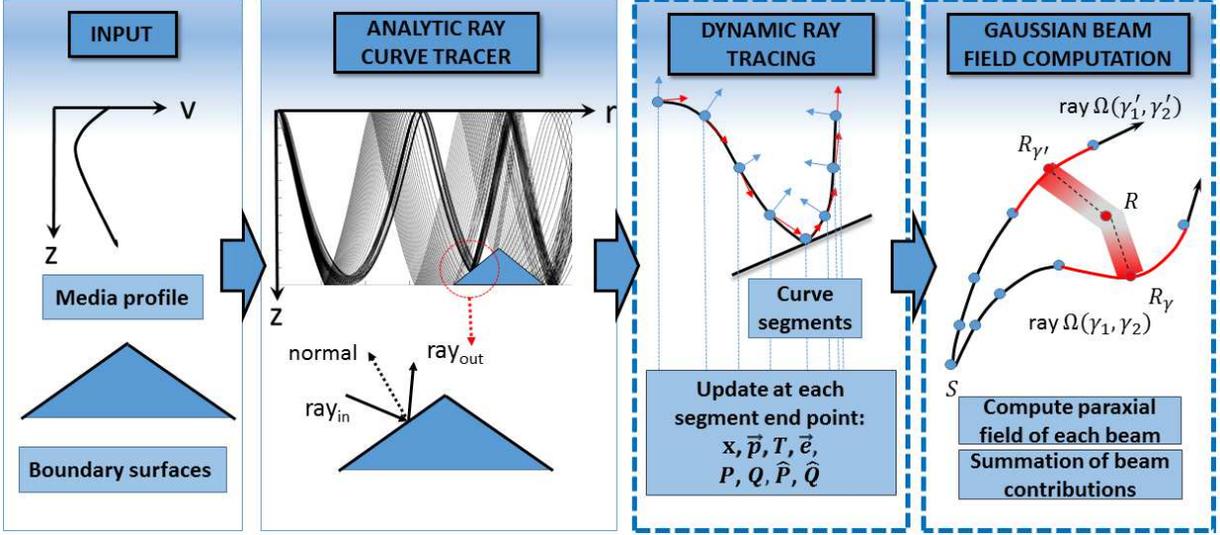


Figure 1: **Algorithm overview.** Given an input media profile and boundary surfaces, the analytic ray tracer⁴⁵ output propagation paths made up of segments of parabolic curves (reviewed in Sec. III A). Our algorithm (the two boxes on the right) then perform dynamic ray tracing that evaluates a set of derivatives analytically and efficiently for pressure along the path (Sec. III B), and the pressure for field point R is computed by summing paraxial contributions from each path based on the Gaussian beam model (Sec. III C). (color online)

evaluations of constant cost, which extends the path-computation efficiency to pressure computation. These variables are subsequently used for computing pressure both on the ray paths (Sec. B) and in the near-ray regions (Sec. C). Here we adopt the term *ray paraxial* defined by Červený⁶ to refer to the near-ray regions, and we use a paraxial Gaussian beam model to approximate the field near a central ray. The mathematical derivation that leads to this algorithm is a special case of the more general ray theory discussed in depth in Červený's comprehensive work⁶; details can be found in the Appendix to this paper.

A. Analytic Ray Curve Tracer

Given a media profile with a spatially-varying sound speed $V(\mathbf{x})$, we assume a locally constant gradient of V^{-2} : $V(\mathbf{x})^{-2} = A_0 + \vec{\mathbf{A}} \cdot \mathbf{x}$ within a certain range around \mathbf{x} . The trajectory of ray originating from \mathbf{x} with initial direction \vec{t}_0 can be shown to be a parabolic curve that lies in the plane with the normal of $\vec{\mathbf{A}} \times \vec{t}_0$ (the *ray plane*, see Figure 2(a)). The intersection between the parabolic ray curve and any planar surface can be solved analytically, and key properties such as position \mathbf{x} , tangent direction \vec{t} , the travel time T , and the *slowness vector* $\vec{p} = \nabla T$ (the direction of which coincides \vec{t}) can be computed for any point along the ray by analytic evaluations of constant cost (Appendix I).

This analytic ray formulation enables a ray tracer that computes propagation paths in a general medium consisting of consecutive segments of parabolic curves. In spatially coherent media, the assumption of constant ∇V^{-2} generally holds for a range larger than the assumption of constant V , enabling the ray curve tracer to advance in longer segments than rectilinear ray tracer; this is one of the key sources of performance improvement. Adaptive segment sizes based on on-the-fly media sampling, as well as acceleration structure that bounds surfaces and ray curves, lead to further speedup that amounts to one to two orders of magnitude improvements over numerical ray integration on 3D scenes⁴⁵.

B. Dynamic Ray Tracing

Dynamic ray tracing, as defined by Červený⁶, keeps track of how a set of derivatives in ray-centered coordinates progress among propagating rays; the derivatives are

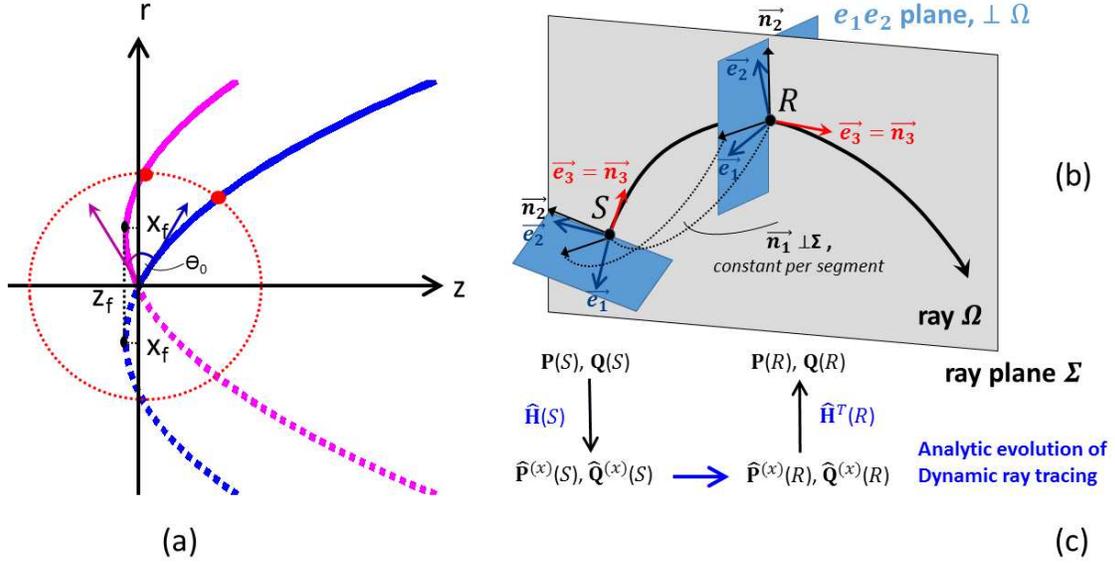


Figure 2: **(a) Analytic ray curve segment** in the ray-plane. Rays of two different initial directions \vec{t}_0 are drawn in blue and magenta. (x_f, z_f) is the vertex of the parabola. We assume a locally-constant ∇V^{-2} , the assumption's range of validity determines the extent of the ray segment. The two red dots mark the end points of ray segments, where the rays leave the range of validity (Sec. A). **(b,c) Analytic dynamic ray tracing:** analytic evolution of \mathbf{P} , \mathbf{Q} are performed for each segment by transforming into and evolving in Cartesian coordinates before transforming back to ray-centered coordinates (Sec. B). The unit basis $\vec{e}_1, \vec{e}_2, \vec{e}_3$ and $\vec{n}_1, \vec{n}_2, \vec{n}_3$ are defined in Appendix III.(color online)

subsequently used to compute pressure and travel time along the ray. We perform dynamic ray tracing analytically for each segment of parabolic ray curve, which achieves efficiency in pressure computation that matches the efficiency of the analytic ray tracer.

First we define the coordinates involved. The *ray-centered coordinates* q_1, q_2, q_3 are defined for any point along a particular ray Ω with origin at that point. The q_3 -axis follows Ω ; the q_1 -axis and q_2 -axis are taken to be perpendicular to the q_3 -axis as well as being

perpendicular to one another. Transformation from the ray-centered coordinates q_k to Cartesian coordinates x_i is accomplished by the 3×3 matrix \hat{H} , $\hat{H}_{ik} = \partial x_i / \partial q_k$, $i, k = 1, 2, 3$, and $\hat{H}^{-1} = \hat{H}^T$ transforms Cartesian coordinates back to ray-centered coordinates.

The derivatives we seek in dynamic ray tracing capture the changes in spatial relationships among rays traveling through a media profile. Consider a system of rays starting from a source and parameterized by *ray parameters* γ_1, γ_2 , taken here as the azimuth ϕ_0 and elevation i_0 angles. The 2×2 matrices \mathbf{Q} and \mathbf{P} are defined with elements $Q_{IJ} = (\partial q_I / \partial \gamma_J)_{T=const}$, $P_{IJ} = (\partial p_I^{(q)} / \partial \gamma_J)_{T=const}$, $I, J = 1, 2$, which are derivatives of the ray-centered coordinates and the slowness vector in the ray-centered coordinates with respect to the ray parameters. Correspondingly, $\hat{\mathbf{Q}}^{(x)}$ and $\hat{\mathbf{P}}^{(x)}$ are defined with elements $Q_{iJ}^{(x)} = (\partial x_i / \partial \gamma_J)_{\sigma=const}$, $P_{iJ}^{(x)} = (\partial p_i^{(x)} / \partial \gamma_J)_{\sigma=const}$, $i = 1, 2, 3$, $J = 1, 2$ in Cartesian coordinates, and $\hat{\mathbf{Q}}^{(x)} = \hat{\mathbf{H}}\mathbf{Q}$, $\hat{\mathbf{P}}^{(x)} = \hat{\mathbf{H}}\mathbf{P}$. Along a parabolic ray curve, $\hat{\mathbf{Q}}^{(x)}$ and $\hat{\mathbf{P}}^{(x)}$ can be evolved analytically (Appendix II), and \hat{H} can also be evolved analytically (Appendix III); we thereby achieve analytic evolution of \mathbf{P} and \mathbf{Q} ((Appendix IV))(See Fig. 2(b,c)).

The pressure amplitude P^{ray} at any point s along the ray given $P(s_0)$ at source s_0 is:

$$P^{ray}(s) = \left[\frac{\rho(s)V(s)\mathbf{J}(s_0)}{\rho(s_0)V(s_0)\mathbf{J}(s)} \right]^{1/2} P(s_0), \quad (1)$$

where ρ is the density of the media, and $\mathbf{J} = \det \mathbf{Q}$ is the *ray Jacobian*. Both \mathbf{P} and \mathbf{Q} are

also used in computing the paraxial field, as explained in Sec. C.

C. Field Computation with Gaussian Beam

We use Gaussian beam model to approximate paraxial fields in the vicinity of ray paths; this involves computing paraxial travel time and paraxial pressure amplitude, both of which benefit from the efficient dynamic ray tracing described in Sec. B above. As shown in the rightmost block of Fig. 1, we gather all segments of ray curves that pass in the vicinity of a given field point, compute the paraxial pressure amplitude and travel time from each segment, and sum up the contribution (Appendix VI). For a field of large volume, we save the costs of locating vicinity ray segments for each field point by reversing the process, distributing the pressure contribution from each beam to the field points it covers.

Based on the definition of \mathbf{P} and \mathbf{Q} , we introduce the 2×2 matrix

$$\mathbf{M} = \mathbf{P}\mathbf{Q}^{-1}, \quad M_{IJ} = (\partial p_I^{(q)} / \partial q_J)_{T=const}, \quad I, J = 1, 2. \quad (2)$$

Recall that the slowness vector \vec{p} is the first derivative of T ; M is therefore the second derivative of T with respect to ray-centered coordinates. For a point R' in the vicinity of a ray Ω , the paraxial travel time at R' can be computed given the T at a point R on Ω :

$$T(R', R) = T(R) + \frac{1}{2} \mathbf{q}^T(R') \mathbf{M}(R) \mathbf{q}(R'), \quad \mathbf{q} = (q_1, q_2)^T, \quad (3)$$

when Ω^\perp is the plane perpendicular to Ω that passes R' , and point R is the intersection of the ray Ω and Ω^\perp . The derivatives of T can also be approximated in Cartesian coordinates, in which case any point R_γ on the ray that is close to R' can be selected, saving the costs of computing Ω^\perp and R (Appendix V).

The Gaussian beam model computes a paraxial amplitude centered on the ray with a Gaussian drop-off, which is achieved by allowing the matrix \mathbf{M} to be complex:

$\mathbf{M} = \text{Re}(\mathbf{M}) + \text{Im}(\mathbf{M})$. $\text{Im}(\mathbf{M})$ is chosen to be positive definite, so that

$$p^{beam}(R') = P^{ray}(R) \exp[-i\omega(-T(R) - \frac{1}{2}\mathbf{q}^T(R')\text{Re}(\mathbf{M}(R))\mathbf{q}(R'))] \quad (4)$$

$$\times \exp[-\frac{1}{2}\omega\mathbf{q}^T(R')\text{Im}(\mathbf{M}(R))\mathbf{q}(R')]. \quad (5)$$

Matrices with suffix a (\mathbf{M}^a , \mathbf{P}^a , \mathbf{Q}^a) represents the matrices of the *actual* field (Eq. 2).

IV. VALIDATION

We validated our algorithm on two benchmark scenes of atmospheric and oceanic sound propagation. The first benchmark was proposed⁴² with reference results generated by a few different methods that agree with each other, and has since been widely adopted by other atmospheric acoustics works for validation purposes. The second benchmark comes from ocean acoustics, with the representative Munk profile and a conical seamount as bathymetry. Results have been reported for the Munk profile in many works, and one of

the latest works⁴³ contains reference results generated by normal modes. We compute 2D pressure fields for these benchmarks and compare the results directly with published results in the literature. Although the ray-based method is a high-frequency approximation, we achieve good agreement with the reference results, and we were able to replicate the characteristic interference patterns for frequency as low as 10Hz. After establishing the validity of our method with these benchmarks, we demonstrate the application of our method to a realistic 3D scene with more complex media conditions in Section V. For all benchmarks, we compute the sound field in terms of the transmission loss (TL), which is defined as: $TL = 20 \log \frac{(Acoustic\ pressure\ at\ a\ field\ point)}{(Acoustic\ pressure\ of\ free\ field\ at\ 1m\ from\ source)}$.

A. Benchmark A (*Inhomogeneous atmosphere, flat ground with impedance*)

The benchmark for outdoor sound propagation proposed in Attenborough et al.⁴² include multiple cases representing different atmospheric conditions; the results generated by a range of methods (including FFP, PE, normal modes, ray and beam tracing) show good agreement on this benchmark. The boundary surface in the scene is a flat ground with impedance, while the media is inhomogeneous with three different profiles.

Media profile: sound speed $c(x)$ at spatial location x with height $z(x)$ is given by

Case 1: Downward refractive $c(x) = 343 + 0.1 * z(x)(m/s)$;

Case 2: Upward refractive $c(x) = 343 - 0.1 * z(x)(m/s)$;

Case 3: Duct condition: case 1 for $z(x) < 100m$, case 2 for $100m < z(x) < 300m$,

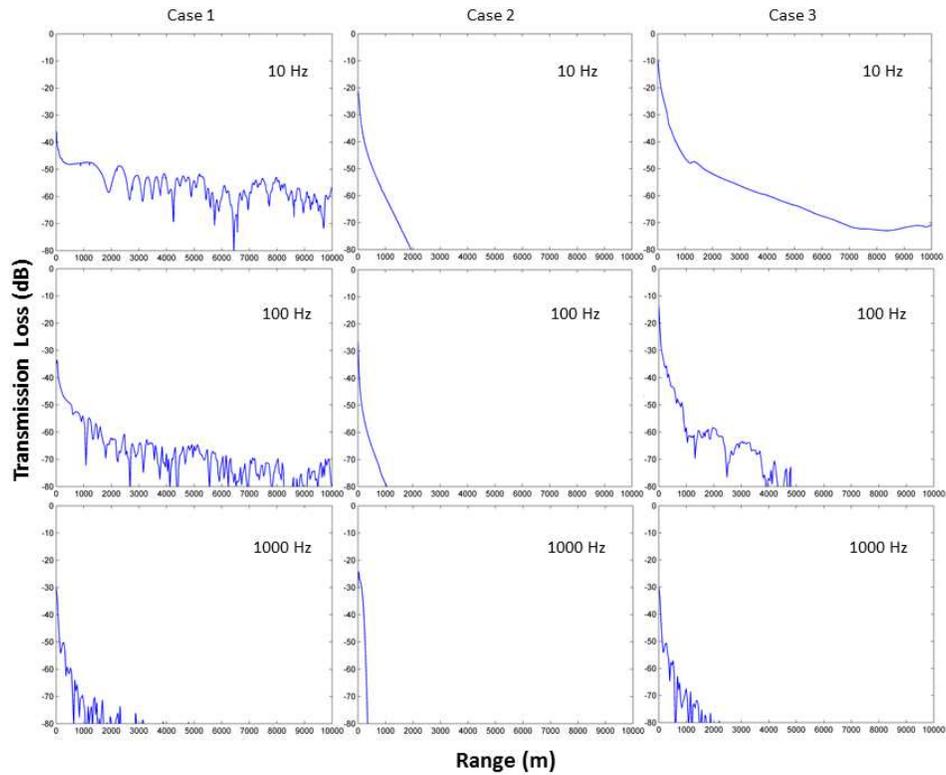


Figure 3: **Benchmark A Range-TL Plot:** Source height $h_s = 5m$, receiver height $h_r = 1m$, range $10km$. Columns contain results for media profiles case 1, 2, and 3, and rows for frequency 10, 100, and 1000 Hz. Direct comparison to Fig. 12-14⁴².(color online)

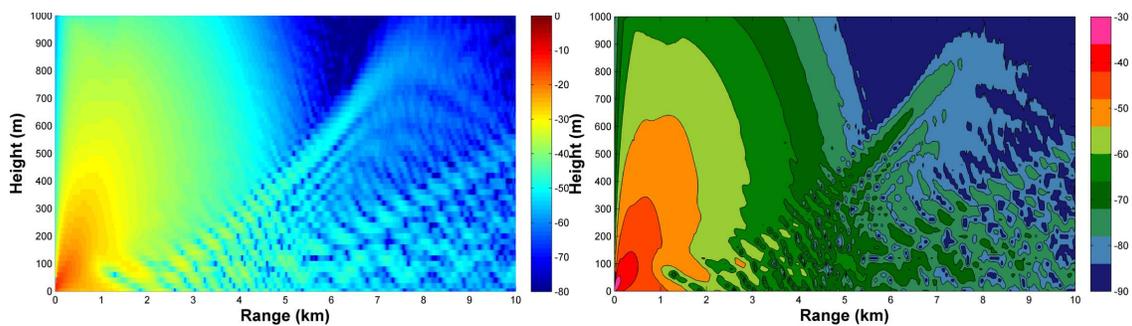


Figure 4: **Benchmark A Field:** Source height $h_s = 5m$, receiver height $h_r = 1m$, frequency 10Hz. The 2D field of height up to $1km$ and range up to $10km$ is visualized on the left. The corresponding contour plot is shown on the right for direct comparison with Fig. 15⁴².(color online)

and constant c for $z(x) > 300m$.

Ground impedance: We use a four-parameter model (with the same parameters from Attenborough et al.⁴²) to compute the impedance of the flat ground.

Results: As shown in Fig. 3 and 4, our algorithm is able to replicate the results for all three cases at three different frequencies. The 1D TL-range plot and the 2D vertical field is shown, which can be compared directly with corresponding figures⁴² listed in the captions. The number of rays required to achieve the results is as low as 21 rays, and more rays can be traced to compute asymptotically more accurate pressure fields.

B. Benchmark B (*Munk profile with conical seamount*)

We validate our algorithm on an underwater benchmark with the standard Munk profile and a conical seamount as bathymetry. The Munk profile is an idealized profile that describes the sound speed variation for depth up to $5000m$. This benchmark spans a much larger range and depth than the atmospheric benchmark above, has a non-linear sound speed variation, and contains a seamount obstacle in the scene. Published results for this benchmark, computed by normal modes, can be found in prior work⁴³.

Media profile: $c(z) = 1500[1 + \epsilon(\tilde{z} - 1 + e^{-\tilde{z}})]$, $\epsilon = 0.00737$, $\tilde{z} = \frac{2(z-1300)}{1300}$.

Bathymetry: conical seamount at range $100km$, height $1000m$, radius of base $20km$.

Bottom impedance: fluid half space with compressional speed of $2000m/s$, density of $1g/cm^3$, and attenuation of $0.1dB/\lambda$.

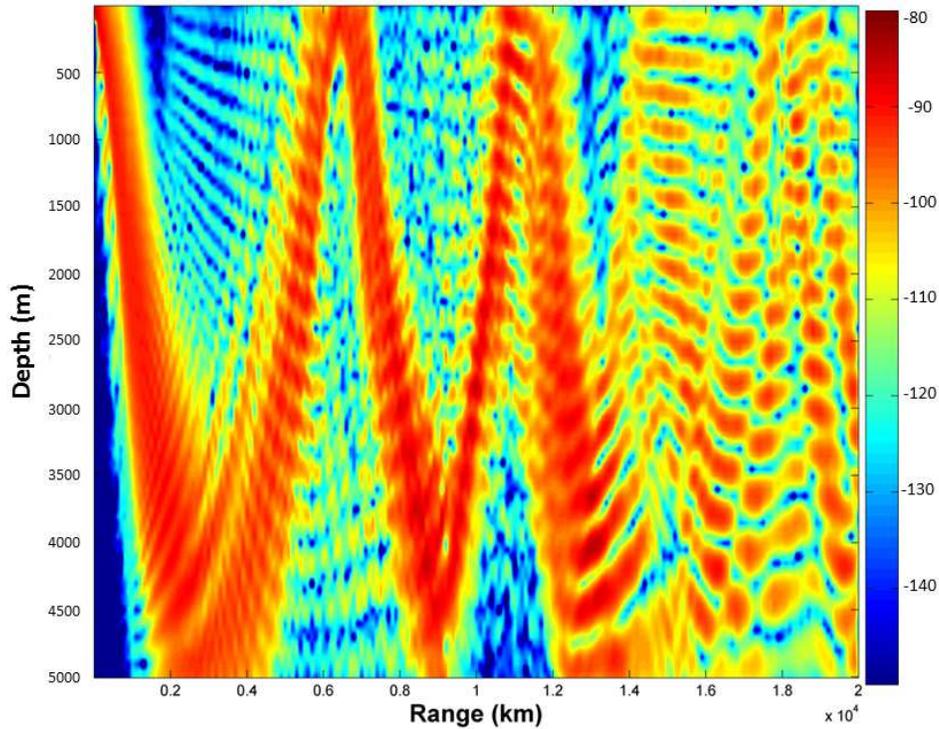


Figure 5: **Benchmark B field:** Source depth 1000m, frequency 50Hz. 2D vertical field is visualized for depth up to 5km and range up to 200km, compared to Fig. 9⁴³.(color online)

Results: The vertical 2D field is visualized in Fig. 5. For this underwater benchmark with extensive volume, our ray tracer successfully replicated the characteristics of the field; compare with the published results computed by normal modes available in prior work⁴³.

V. APPLICATION ON COMPLEX OUTDOOR SCENE

A. Scene configuration

Media We generate a general media profile based on a commonly-used empirical model of the atmosphere¹. The acoustic index of refraction in the atmosphere ($n = c_0/c$, where c_0 is

the reference sound speed) is modeled with a stratified component n_{str} and a fluctuation component n_{flu} , so that $n = n_{str} + n_{flu}$. The stratified component follows a logarithmic profile of the altitude z : $n_{str}(z) = c_0 / (c_0 + b \ln(\frac{z}{z_g} + 1))$, where z_g is the roughness length of the ground surface, and a typical value for b is 1 m/s for a downward-refracting atmosphere and -1 m/s for an upward-refracting atmosphere. The fluctuation component at position \mathbf{x} can be computed as $n_{flu}(\mathbf{x}) = \sum_i G(\mathbf{k}_i) \cos(\mathbf{k}_i \cdot \mathbf{x} + \varphi_i)$, where \mathbf{k}_i is the wave vector describing the *spatial* frequency of the fluctuation, φ_i is a random angle between $[0, 2\pi]$, and $G(\mathbf{k}_i)$ is a normalization factor. This model represents typical atmospheric conditions with a realistic level of variations. Profiles from empirical measurements or other functions and models can be substituted as input to our algorithm.

Scene Objects We use a computer-modeled 3D scene consisting of undulating terrains with a reservoir and buildings. A wireframe rendering of the scene, with the two sound-source locations marked by green and red dots, can be found in Fig. 6(a). The scene has a physical dimension of $220m \times 150m \times 50m$, and is represented by 4,000 triangular surface primitives. Our algorithm can simulate propagation for any scenes that can be modeled or scanned into surface representations similar to the one demonstrated here.

B. Results

We compute sound fields for two source locations (Fig. 6(a)), one on the slope of the reservoir, the other in the center of the valley. For this scene, the sound field displays

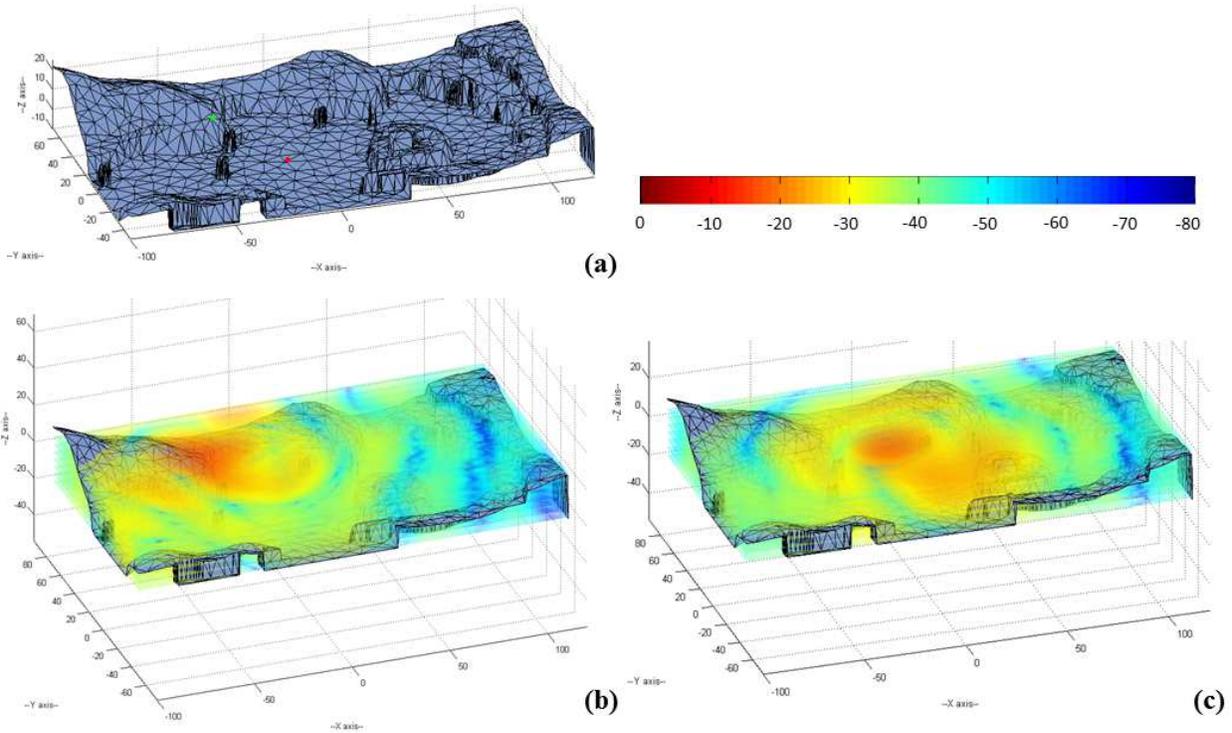


Figure 6: **Outdoor scene and field results** for two source locations: (a) Wireframe rendering of the Reservoir scene. The green dot represents a sound source located on the slope, and the red dot represents a sound source in the valley. (b) Slices of sound TL field visualized for the green source. (c) Slices of sound TL level visualized for the red source. Frequency 10Hz.(color online)

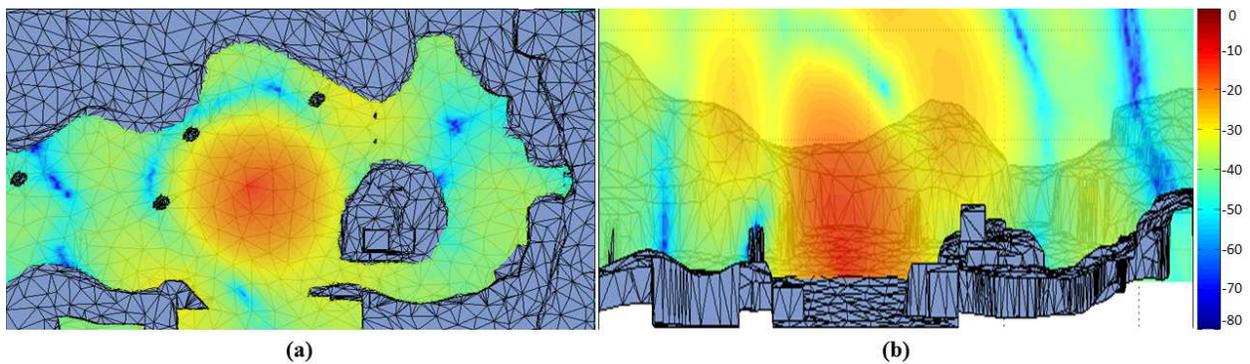


Figure 7: **Source in the valley: Downward-refractive media.** (a) horizontal slice and (b) vertical slice, both of which pass the source location. Frequency 10Hz.(color online)

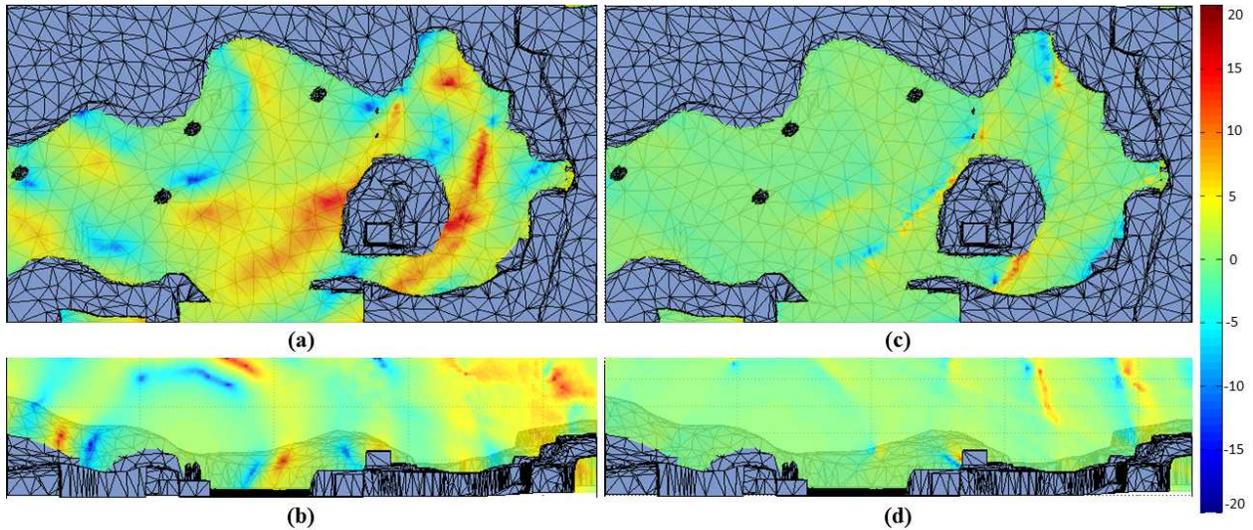


Figure 8: **Source on the slope: (a) upward vs. downward refraction** relative sound pressure level, top view of a horizontal slice of the field; **(b) front view of a vertical slice of the same field as (a).** **(c) up wind vs. down wind** relative sound pressure level, top view of a horizontal slice of the field; **(d) front view of a vertical slice of the same field as (c).** Frequency 10Hz.(color online)

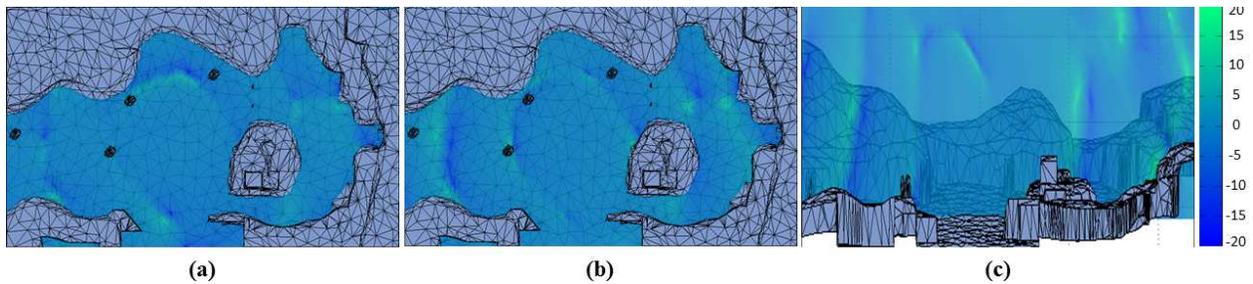


Figure 9: **Vector wind for source in the valley: (a) Difference in horizontal field of sound pressure level between a north and a south wind.** **(b) Difference in horizontal field of sound pressure level between an east and a west wind.** **(c) Difference in vertical field of sound pressure level between an east and a west wind.** Frequency 10Hz.(color online)

characteristics resulting from the interaction of sound waves with the medium itself and with obstacles in the scene. We show a stack of slices in Fig. 6 for visualization purposes, while the full 3D field is computed by our algorithm. A source frequency of 10Hz is used for all the visualized results in order to keep the field pattern easy to interpret. Our algorithm generate the sound fields at close to interactive rates (679 *ms/frame* tracing 1K initial rays for the $220 \times 150 \times 50$ field points). At this rate, insights about field characteristics can be gained by varying the scene configurations and observe the field changes. Details on performance of our method can be found by referring to⁴⁵.

The diurnal changes in the atmosphere typically lead to an upward-refractive condition during the day and a downward-refractive condition at night. For the source in the valley, we show a horizontal and a vertical slice of the field under a downward-refractive condition in Fig. 7. For the source on the slope, we visualize the relative sound pressure level between an upward and a downward refractive conditions, again showing a horizontal and a vertical slice in Fig. 8(a) and (b), respectively.

Wind plays an important role in atmospheric sound propagation, creating extra variations in the media profile and interacting with obstacles that further complicate the sound field. For the sound source on the slope, we simulate the sound field for up-wind and down-wind conditions (Fig. 8(c) and (d)); this yields patterns similar (in their relative difference) to the difference between upward- and downward- refractive media. Our ray

tracer also accounts for media with a vector wind field (see⁴⁵). For the source in the valley, we show this capability by visualizing the differences in the fields between a west and east wind (Fig. 9(a)), and between a north and south wind (Fig. 9 (b,c)).

VI. DISCUSSIONS

This algorithm is complementary to many existing sound propagation techniques and can be extended or combined in multiple ways. With the analytic ray curve tracer⁴⁵ as a component, our method inherits its many potential extensions, such as: augmenting GA methods with capability to handle inhomogeneous media; forming hybrid method with ESM based on frequency and spatial subdivision (similar to Yeh et al.²⁷); using the ray tracer for wide area assessment that guides the application of numerical methods only in areas of interest. As in other ray models, extensions^{38;39} to our method can also be built to account for turbulence. It is also possible to accommodate sound sources other than a point source, such as complex sources or sources with directivity^{40;41}; this can be achieved by using the techniques of Gaussian beam expansion.

As a ray-based model, this algorithm inherits the limitations that it is a high-frequency approximation, not a full-wave solution. The analytic ray tracer relies on spatial coherence in the medium to perform efficiently. The Gaussian beam model that is used to compute the sound field relies on carefully chosen parameters that control the beam width³⁷, and it is best determined on a per-scene basis.

VII. CONCLUSION

This paper combines an analytic ray curve tracer⁴⁵ and the Gaussian beam model to form an efficient solution for outdoor sound field computation. Based on the parabolic ray formulation⁴⁵, we use analytic solutions to compute on-ray pressure and paraxial fields in combination with a Gaussian beam model. The analytic ray tracer's path-computation efficiency is matched by our pressure-computation efficiency, and the combined algorithm can simulate the propagated sound field for large three-dimensional outdoor scenes with general input media and complex obstacles. This algorithm is validated on 2D benchmarks with inhomogeneous media profiles that are widely used in atmospheric and underwater propagation. The results computed by our algorithm are verified against published results, generated by validated methods including FDTD, PE, and normal modes. The capability of this algorithm is further demonstrated with a complex 3D scene under a variety of media conditions that would present scalability challenges to existing methods. Results that reflect the characteristics of the scene and media are generated at close-to-interactive speed. As future work we hope to obtain measured data or to run large scale numerical simulation to further validate and gauge the speedup of the algorithm, and to apply this algorithm on more challenging outdoor scenarios, including fully dynamic scenes.

APPENDIX

I. Analytic evolution of ray trajectories

With ray parameter σ defined by $d\sigma = V^2 dT = V ds$, the ray trajectories x_i , slowness p_i , and travel time T can be evolved analytically from σ_0 to any σ along the ray:

$$x_i(\sigma) = x_{i0} + p_{i0}(\sigma - \sigma_0) + \frac{1}{4}A_i(\sigma - \sigma_0)^2, \quad (6)$$

$$p_i(\sigma) = p_{i0} + \frac{1}{2}A_i(\sigma - \sigma_0), \quad (7)$$

$$T(\sigma) = T(\sigma_0) + V_0^{-2}(\sigma - \sigma_0) + \frac{1}{2}A_i p_{i0}(\sigma - \sigma_0)^2 + \frac{1}{12}A_i A_i(\sigma - \sigma_0)^3. \quad (8)$$

II. Analytic evolution of Cartesian P and Q

The characteristic system of the Hamiltonian for of the Eikonal equation gives:

$$\frac{dp_i}{d\sigma} = \frac{1}{2} \frac{\partial}{\partial x_i} \left(\frac{1}{V^2} \right), \quad \frac{dT}{d\sigma} = p_k p_k = V^{-2}. \quad (9)$$

From equation 9 and because partial derivative $\partial/\partial\gamma$ commutes with $d/d\sigma$, a simple *dynamic ray tracing system* can be derived as follows:

$$\frac{d}{d\sigma} Q_i^{(x)} = P_i^{(x)}, \quad \frac{d}{d\sigma} P_i^{(x)} = \frac{1}{2} \frac{\partial^2}{\partial x_i \partial x_j} \left(\frac{1}{V^2} \right) Q_j^{(x)}, \quad (10)$$

For constant media gradient of V^{-2} , 10 can be solved analytically for any point R at σ along the ray Ω if $\hat{\mathbf{Q}}^{(x)}$ and $\hat{\mathbf{P}}^{(x)}$ are known at any other point S at σ_0 along the ray Ω :

$$P_{iJ}^{(x)}(x_i) = P_{iJ}^{(x)}(S), \quad Q_{iJ}^{(x)}(R) = Q_{iJ}^{(x)}(S) + (\sigma - \sigma_0)P_{iJ}^{(x)}(S). \quad (11)$$

III. Analytic evolution of transformation matrix

For constant gradient of V^{-2} , \hat{H} can be solved analytically for any point R from any other point S along Ω . This is achieved by computing the ray-centered coordinates unit basis \vec{e}_1 , \vec{e}_2 , and \vec{e}_3 that constitutes \hat{H} , as $\hat{H}_{ik} = \partial x_i / \partial q_k = \partial q_k / \partial x_i = e_{ki}$, $i, k = 1, 2, 3$. Consider a set of orthonormal unit vectors \vec{n}_1 , \vec{n}_2 , \vec{n}_3 defined along ray Ω , parameterized by σ . Let $\vec{n}_3(\sigma) = V(\sigma)\vec{p}(\sigma)$ follow the tangent of the ray, $\vec{n}_2(\sigma)$ is selected to be perpendicular to the ray plane, \vec{n}_1 is then defined by $\vec{n}_1 = \vec{n}_2 \times \vec{n}_3$. Because the ray is a planar curve for constant gradient V^{-2} , $\vec{n}_1(\sigma) = \vec{n}_1(\sigma_0)$. Given Equations 6-8,

$$\begin{aligned} \vec{n}_2(\sigma) &= \vec{n}_1(\sigma) \times \vec{n}_3(\sigma) = \vec{n}_1(\sigma) \times V(\sigma)\vec{p}(\sigma) \\ &= \vec{n}_1(\sigma_0) \times V(\sigma)(\vec{p}(\sigma_0) + \frac{1}{2}\vec{A}(\sigma - \sigma_0)), \end{aligned} \quad (12)$$

As e_3 coincides with \vec{n}_3 , $\vec{e}_1(\sigma)$, $\vec{e}_2(\sigma)$ can be determined from $\vec{e}_1(\sigma_0)$, $\vec{e}_2(\sigma_0)$ and the

evolution of \vec{n}_1, \vec{n}_2 from σ_0 to σ is:

$$\begin{aligned}\vec{e}_1(\sigma) &= [\vec{e}_1(\sigma_0) \cdot \vec{n}_1(\sigma_0)]\vec{n}_1(\sigma) + [\vec{e}_1(\sigma_0) \cdot \vec{n}_2(\sigma_0)]\vec{n}_2(\sigma), \\ \vec{e}_2(\sigma) &= [\vec{e}_2(\sigma_0) \cdot \vec{n}_1(\sigma_0)]\vec{n}_1(\sigma) + [\vec{e}_2(\sigma_0) \cdot \vec{n}_2(\sigma_0)]\vec{n}_2(\sigma).\end{aligned}\tag{13}$$

IV. Evolution of ray-centered \mathbf{P} and \mathbf{Q}

1. Take initial condition for \mathbf{P}, \mathbf{Q} . Assuming a point source S and ϕ_0 and i_0 as the ray

parameters γ_1, γ_2 : $\mathbf{Q}(S) = 0, \quad \mathbf{P}(S) = \frac{1}{v(S)} \begin{pmatrix} 1 & 0 \\ 0 & \sin i_0 \end{pmatrix}.$

2. Transform \mathbf{P}, \mathbf{Q} into $\hat{\mathbf{P}}^{(x)}$ and $\hat{\mathbf{Q}}^{(x)}$ with \hat{H} ,
3. Analytically evolve $\hat{\mathbf{P}}^{(x)}$ and $\hat{\mathbf{Q}}^{(x)}$ by Equation 11,
4. Analytically evolve $\hat{\mathbf{H}}$ by Equation 13,
5. Transform the evolved $\hat{P}^{(x)}$ and $\hat{Q}^{(x)}$ back to \mathbf{P}, \mathbf{Q} with the evolved \hat{H}^T .

V. Cartesian paraxial travel time

Denote the Cartesian coordinates of R' and R_γ by $x_i(R')$ and $x_i(R_\gamma)$, and $x_i(R, R_\gamma) = x_i(R) - x_i(R_\gamma)$, the quadratic expansion of T from $T(R_\gamma)$ is:

$$T(R, R_\gamma) = T(R_\gamma) + x_i(R, R_\gamma)p_i^{(x)}(R_\gamma) + \frac{1}{2}x_i(R, R_\gamma)x_j(R, R_\gamma)M_{ij}^{(x)}(R_\gamma),\tag{14}$$

where M_{ij} are the elements of the 3×3 matrix $\hat{\mathbf{M}}^{(x)}$:

$$\hat{\mathbf{M}}(R_\gamma) = \hat{\mathbf{H}}(R_\gamma) \begin{pmatrix} \mathbf{M}(R_\gamma) & M_{13}(R_\gamma) \\ & M_{23}(R_\gamma) \\ M_{13}(R_\gamma) & M_{23}(R_\gamma) & M_{33}(R_\gamma) \end{pmatrix} \hat{\mathbf{H}}^T(R_\gamma). \quad (15)$$

Here $\mathbf{M}(R_\gamma)$ is defined in Eq. 2. The remaining elements can be derived⁶ to be:

$$M_{13}(R_\gamma) = -(v^{-2}v_{,1})_{R_\gamma}, \quad M_{23}(R_\gamma) = -(v^{-2}v_{,2})_{R_\gamma}, \quad M_{33}(R_\gamma) = -(v^{-2}v_{,3})_{R_\gamma}, \quad (16)$$

$$v = [V(q_1, q_2, s)]_{q_1=q_2=0, s=s(R_\gamma)}, \quad v_{,i} = [\partial V(q_1, q_2, s)/\partial q_i]_{q_1=q_2=0, s=s(R_\gamma)}. \quad (17)$$

$v_{,i}$ can be solved by transforming to Cartesian coordinates first: $v_{,i} = \partial V/\partial q_i = H_{ki}\partial V/\partial x_k$, and $\partial V/\partial x_k$ can be solved analytically for constant gradient of V^{-2} by:

$$\partial V^{-2}/\partial x_k = -2V^{-3}\partial V/\partial x_k = A_k \implies \partial V/\partial x_k = -\frac{1}{2}V^3 A_k. \quad (18)$$

VI. Gaussian beam summation

The contributions of Gaussian beams are then summed up by integral superposition:

$$p(R, \omega) = \iint_{\mathcal{D}} \Phi(\gamma_1, \gamma_2) P^{ray}(R_\gamma) \exp[i\omega T(R, R_\gamma)] d\gamma_1 d\gamma_2, \quad (19)$$

where R is the field point and R_γ is a point on the ray γ of the ray parameter γ_1, γ_2 . The weighting function Φ is derived to be:

$$\Phi(\gamma_1, \gamma_2) = (\omega/2\pi)[- \det(\mathbf{M}(R_\gamma) - \mathbf{M}^a(R_\gamma))]^{1/2} |\det \mathbf{Q}^a(R_\gamma)| \quad (20)$$

$$= (\omega/2\pi)[- \det(\mathbf{Q}^{aT}(\mathbf{M} - \mathbf{M}^a)\mathbf{Q}^a)]^{1/2}. \quad (21)$$

The choice of $\text{Re}(\mathbf{M})$ is related to the curvatures of the wavefront and the choice of $\text{Im}(\mathbf{M})$ is related to the width of the amplitude profile. They can be specified at R_γ or any other point along the central ray γ to control the shape of the beam.

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